

電磁場是向量還是張量？

電磁場到底是向量還是張量？答案可說「以上皆是」，亦可說是「以上皆非」。聽來雖然荒謬，但我們能否從中得到什麼啟示(如「假向量」和「反夸克」的本質)？

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最簡單的問題**vs.**最簡單的答案

- 當然是張量！

錯誤的問題vs.錯誤的答案

- 錯誤的答案
- 錯誤的問題

該怎麼問才算正確？

- 什麼是張量？
- 方陣＝二階張量？
- M_{ij} ：二階張量的分量？
- $1 \times n$ 矩陣＝向量（一階張量）？
- V_i ：向量的分量？

張量必須用「變換式」定義

$\mathbb{R}^n$: x -system $\Rightarrow$ y -system

$$V^a(y) = \frac{\partial y^a}{\partial x^m} V^m(x), \text{ contravariant} \quad \text{逆變(反變)}$$

$$V_a(y) = \frac{\partial x^m}{\partial y^a} V_m(x), \text{ covariant} \quad \text{協變(共變)}$$

張量必須用「變換式」定義(續)

$$T^{ab}(y) = \frac{\partial y^a}{\partial x^m} \frac{\partial y^b}{\partial x^n} T^{mn}(x), \begin{pmatrix} 2 \\ 0 \end{pmatrix} \text{ type}$$

$$T_{ab}(y) = \frac{\partial x^m}{\partial y^a} \frac{\partial x^n}{\partial y^b} T_{mn}(x), \begin{pmatrix} 0 \\ 2 \end{pmatrix} \text{ type}$$

$$T^a_b(y) = \frac{\partial y^a}{\partial x^n} \frac{\partial x^n}{\partial y^b} T^n_m(x), \begin{pmatrix} 1 \\ 1 \end{pmatrix} \text{ type.}$$

Linear Transformation

$$y^a = R^a_{\cdot m} x^m; \quad x^m = (R^{-1})^m_{\cdot a} y^a$$

$$V^a(y) = R^a_{\cdot m} V^m(x)$$

$$V' = RV$$

$$V_a(y) = (R^{-1})^m_{\cdot a} V_m(x) = (\tilde{R}^{-1})^{\cdot m}_a V_m(x)$$

Linear Transformation (Cont.)

$$T^{ab}(y) = R^a_{.m} R^b_{.n} T^{mn}(x)$$

$$T_{ab}(y) = (\tilde{R}^{-1})^{.m}_a (\tilde{R}^{-1})^{.n}_b T_{mn}(x)$$

$$T^a_b(y) = R^a_{.m} (\tilde{R}^{-1})^{.n}_b T^m_n(x)$$

Electromagnetic Tensor in Minkowski Space $\mathbb{R}^{3,1}$

$$F_{ab} = \begin{pmatrix} 0 & E_x & E_y & E_z \\ -E_x & 0 & -B_z & B_y \\ -E_y & B_z & 0 & -B_x \\ -E_z & -B_y & B_x & 0 \end{pmatrix}$$

$$F^{ab} = \begin{pmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & -B_z & B_y \\ E_y & B_z & 0 & -B_x \\ E_z & -B_y & B_x & 0 \end{pmatrix}$$

$$F^{ab} = \begin{pmatrix} 0 & -\tilde{\mathbf{E}} \\ \mathbf{E} & \mathbf{M} \end{pmatrix}$$

Electromagnetic Tensor

- Rank-2 tensor under $SO(3,1)$, i.e.
Rank-2 tensor under Lorentz transformation.
- E_i, B_j are components of the rank-2 tensor.

$$F^{ab}(y) = L^a_{\cdot m} L^b_{\cdot n} F^{mn}(x)$$

$$F_{ab}(y) = (\tilde{L}^{-1})^{\cdot m}_a (\tilde{L}^{-1})^{\cdot n}_b F_{mn}(x)$$

$$\begin{aligned}\partial_a F^{ab} &= 4\pi J^b \\ \partial_{[a} F_{bc]} &= 0\end{aligned}$$

$$F^{ab}(y) = L^a_{.m} L^b_{.n} F^{mn}(x)$$

$$\longleftrightarrow F' = LF\tilde{L}$$

$$F^{ab}(y) = L^a_{.m} L^b_{.n} F^{mn}(x)$$

$$= L^a_{.m} F^{mn}(x) L^b_{.n}$$

$$= L^a_{.m} F^{mn}(x) \tilde{L}^{.b}_n$$

$$= [LF(x)\tilde{L}]^{ab}$$

Special case: $SO(3,1) \rightarrow SO(3)$

$$L(SO(3)) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & * & * & * \\ 0 & * & * & * \\ 0 & * & * & * \end{pmatrix} = \begin{pmatrix} 1 & \mathbf{0} \\ \mathbf{0} & R \end{pmatrix}$$

$$LF\tilde{L} = \begin{pmatrix} 1 & \mathbf{0} \\ \mathbf{0} & R \end{pmatrix} \begin{pmatrix} 0 & -\tilde{\mathbf{E}} \\ \mathbf{E} & \mathbf{M} \end{pmatrix} \begin{pmatrix} 1 & \mathbf{0} \\ \mathbf{0} & \tilde{R} \end{pmatrix}$$

$$\begin{aligned}
LF\tilde{L} &= \begin{pmatrix} 1 & \mathbf{0} \\ \mathbf{0} & R \end{pmatrix} \begin{pmatrix} 0 & -\tilde{\mathbf{E}} \\ \mathbf{E} & \mathbf{M} \end{pmatrix} \begin{pmatrix} 1 & \mathbf{0} \\ \mathbf{0} & \tilde{R} \end{pmatrix} \\
&= \begin{pmatrix} 0 & -\tilde{\mathbf{E}} \\ R\mathbf{E} & R\mathbf{M} \end{pmatrix} \begin{pmatrix} 1 & \mathbf{0} \\ \mathbf{0} & \tilde{R} \end{pmatrix} \\
&= \begin{pmatrix} 0 & -\tilde{\mathbf{E}}\tilde{R} \\ R\mathbf{E} & R\mathbf{M}\tilde{R} \end{pmatrix} \\
&= \begin{pmatrix} 0 & -\tilde{\mathbf{E}}' \\ \mathbf{E}' & \mathbf{M}' \end{pmatrix}
\end{aligned}$$

$$\mathbf{E}' = R\mathbf{E}$$

$$\mathbf{M}' = R\mathbf{M}\tilde{R}$$

Under SO(3)

$$\mathbf{E}' = R\mathbf{E} \quad \text{Vector}$$

$$\mathbf{M}' = R\mathbf{M}\tilde{R} \quad \text{Rank-2 Tensor}$$

$$\mathbf{M}^{ij} = \begin{pmatrix} 0 & -B_z & B_y \\ B_z & 0 & -B_x \\ -B_y & B_x & 0 \end{pmatrix}$$

$$\mathbf{M}^{ij} = \begin{pmatrix} 0 & -B_z & B_y \\ B_z & 0 & -B_x \\ -B_y & B_x & 0 \end{pmatrix} \equiv \begin{pmatrix} 0 & -B_3 & B_2 \\ B_3 & 0 & -B_1 \\ -B_2 & B_1 & 0 \end{pmatrix}$$

$$\mathbf{M}^{ij} = -\epsilon^{ijk} B_k, \quad B_i = -\frac{1}{2} \epsilon_{ijk} \mathbf{M}^{jk}$$

Tensor contraction?

Is ϵ_{ijk} a rank-3 tensor under $SO(3)$?

Claim1: ϵ_{ij} is a rank-2 tensor under $SO(2)$

$$\det R = R^1_{.1}R^2_{.2} - R^1_{.2}R^2_{.1} = \epsilon^{mn}R^1_{.m}R^2_{.n} = \frac{1}{2!}\epsilon_{ij}\epsilon^{mn}R^i_{.m}R^j_{.n} = 1$$

$$\epsilon^{mn}\epsilon_{pq} = \delta_p^m\delta_q^n - \delta_q^m\delta_p^n$$

$$\frac{1}{2!}\epsilon_{ij}\epsilon^{mn}\epsilon_{pq}R^i_{.m}R^j_{.n} = \epsilon_{pq}$$

$$\frac{1}{2!}\epsilon_{ij}(\delta_p^m\delta_q^n - \delta_q^m\delta_p^n)R^i_{.m}R^j_{.n} = \epsilon_{pq}$$

$$\frac{1}{2!}\epsilon_{ij}(R^i_{.p}R^j_{.q} - R^i_{.q}R^j_{.p}) = \epsilon_{pq}$$

$$R^i_{.p}R^j_{.q}\epsilon_{ij} = \epsilon_{pq}$$

Claim2: ϵ_{ijk} is a rank-3 tensor under $SO(3)$

$$\det R = \frac{1}{3!} \epsilon_{ijk} \epsilon^{mnl} R^i_{.m} R^j_{.n} R^k_{.\ell} = 1$$

$$R^i_{.p} R^j_{.q} R^k_{.r} \epsilon_{ijk} = \epsilon_{pqr}$$

Constant Tensor

$$R^i_{.m} R^j_{.n} R^k_{.\ell} \epsilon_{ijk} = \epsilon_{mnl}$$

Conclusion:

Under SO(3)

$$1. B_i \in \begin{pmatrix} 0 & -B_3 & B_2 \\ B_3 & 0 & -B_1 \\ -B_2 & B_1 & 0 \end{pmatrix}$$

is a component of a rank-2 tensor.

$$\mathbf{M}^{ij} = -\epsilon^{ijk} B_k, \quad B_i = -\frac{1}{2} \epsilon_{ijk} \mathbf{M}^{jk}$$

2. $B_i \in (B_1, B_2, B_3)$ is a component of a vector.

Case study1. **B** is a pseudo-vector

A pseudo-vector (or axial vector) is a quantity that transforms like a vector under a proper rotation, but in three dimensions **gains an additional sign flip under an improper rotation** such as a reflection.

<http://en.wikipedia.org/wiki/Pseudovector>

$$V' = (\det R)RV = \pm RV$$

$$SO(3) \implies O(3) : R \in O(3), \det R = \pm 1$$

For improper rotation, $\det R = -1$,

$$\det R = \frac{1}{3!} \epsilon_{ijk} \epsilon^{mnl} R^i_{.m} R^j_{.n} R^k_{.\ell} = -1$$

$$\epsilon_{ijk} R^i_{.m} R^j_{.n} = -\epsilon_{mnl} (R^{-1})^{\ell}_{.k} = -\epsilon_{mnl} R^{\ell}_{.k}$$

$$\mathbf{M}^{ij} = -\epsilon^{ijk} B_k, \quad B_i = -\frac{1}{2} \epsilon_{ijk} \mathbf{M}^{jk}$$

ϵ_{ijk} is invariant under transformation.

$$B'_i = -\frac{1}{2} \epsilon_{ijk} \mathbf{M}'^{jk}$$

$$= -\frac{1}{2} \epsilon_{ijk} R^j_{\cdot m} R^k_{\cdot n} \mathbf{M}^{mn}$$

$$= -\frac{1}{2} (-\epsilon_{mnl} R^{\cdot \ell}_i) \mathbf{M}^{mn} \quad \leftarrow \quad \epsilon_{ijk} R^i_{\cdot m} R^j_{\cdot n} = -\epsilon_{mnl} (R^{-1})^{\ell}_{\cdot k} = -\epsilon_{mnl} R^{\cdot \ell}_k$$

$$= -R^{\cdot \ell}_i \left[-\frac{1}{2} \epsilon_{mnl} \mathbf{M}^{mn} \right]$$

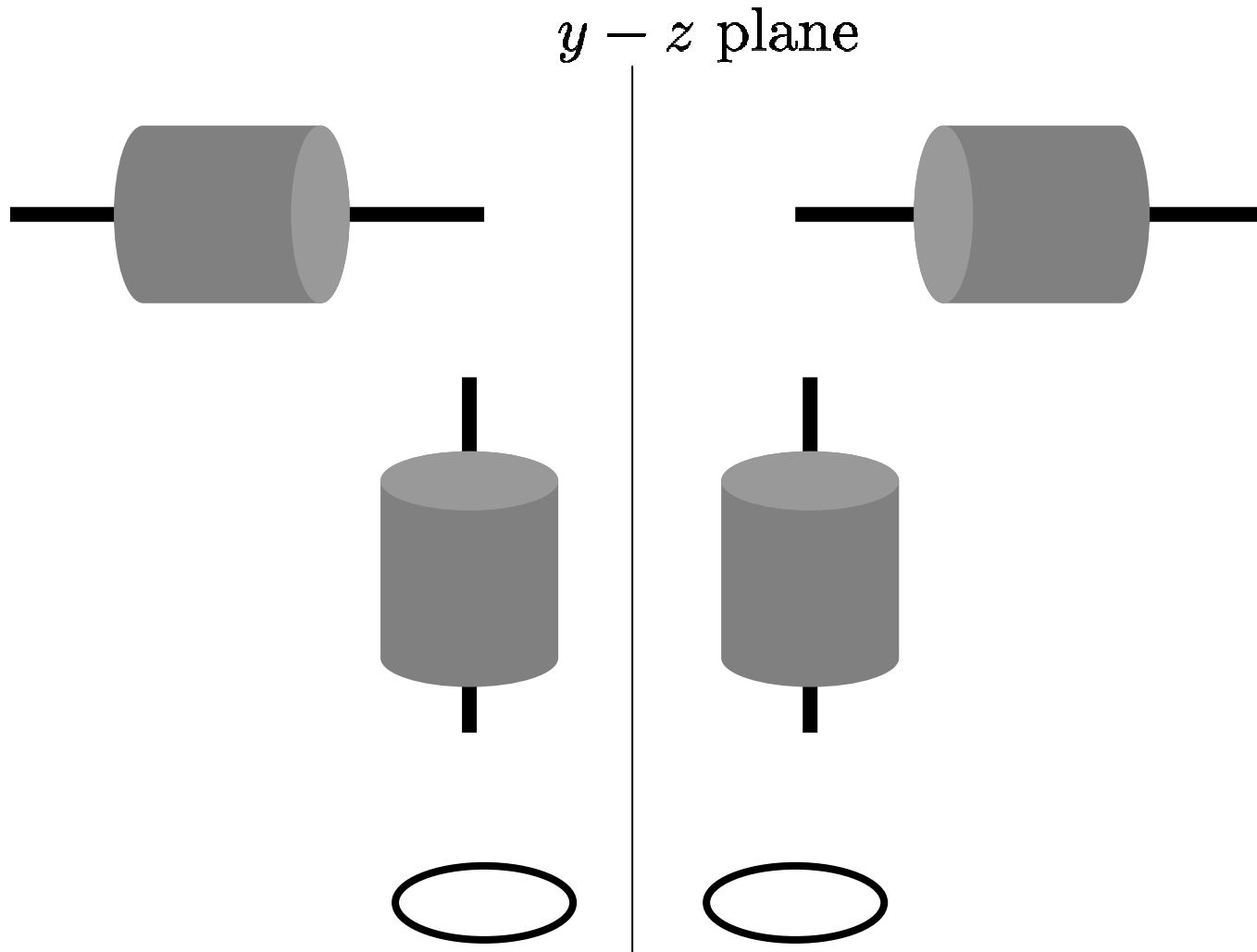
$$= -R^{\cdot \ell}_i B_\ell$$

x -reflection $\mathbf{M}^{ij} = \begin{pmatrix} 0 & -B_z & B_y \\ B_z & 0 & -B_x \\ -B_y & B_x & 0 \end{pmatrix}, \quad \mathbf{M}' = R\mathbf{M}\tilde{R}$

$$\begin{aligned} \mathbf{M}'^{ij} &= \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 0 & -B_z & B_y \\ B_z & 0 & -B_x \\ -B_y & B_x & 0 \end{pmatrix} \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \\ &= \begin{pmatrix} 0 & B_z & -B_y \\ -B_z & 0 & -B_x \\ B_y & B_x & 0 \end{pmatrix} \end{aligned}$$

$$\longleftrightarrow \mathbf{B}' = - \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \mathbf{B} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix} \mathbf{B}$$

Ampère's right hand rule → left hand rule?



Case study 2.

1. Cross product of two vectors $A \times B$ is still a vector.

2. Cross product of two vectors $A \times B$ is a pseudovector.

2'. Cross product of two polar vectors $A \times B$ is an axial vector.

Case study 3.

What's special for anti-symmetry

T^{ab} is a general rank-2 tensor in $\mathbb{R}^3$.

$$T^{ab} = S^{ab} + A^{ab}$$

$$\frac{1}{2!}\epsilon_{abc}T^{ab} = \frac{1}{2!}\epsilon_{abc}(S^{ab} + A^{ab}) = \frac{1}{2!}\epsilon_{abc}A^{ab}$$



Rearranging the component

Case study 4.

Generalization by an example

T^{abc} is a rank-3 tensor under $SO(5)$ in $\mathbb{R}^5$,

Then $\frac{1}{3!}\epsilon_{abcde}T^{abc}$ is a rank-2 tensor.

$$\frac{1}{3!}\epsilon_{abcde}T^{abc} = \frac{1}{3!}\epsilon_{abcde}(S^{abc} + M^{abc} + A^{abc}) = \frac{1}{3!}\epsilon_{abcde}A^{abc}$$

Case study 5. Anti-quarks

Quarks $\in \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ type tensor under $SU(3)$.

$\begin{pmatrix} 2 \\ 0 \end{pmatrix}$ type anti-symmetric tensor $\equiv \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ type tensor.

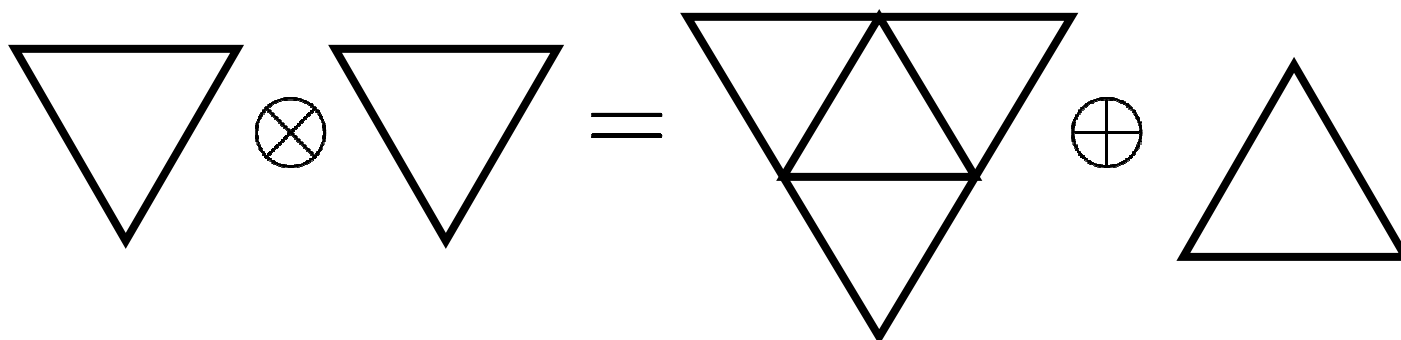
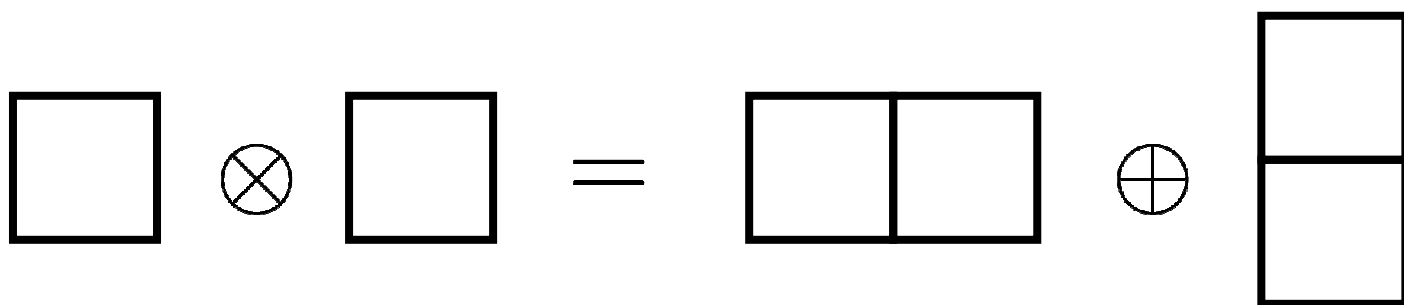
$$ud - du \Rightarrow \bar{s}$$

$$ds - sd \Rightarrow \bar{u}$$

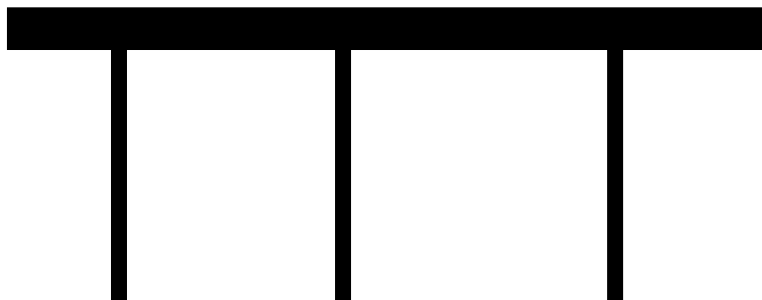
$$su - us \Rightarrow \bar{d}$$

SU(3)

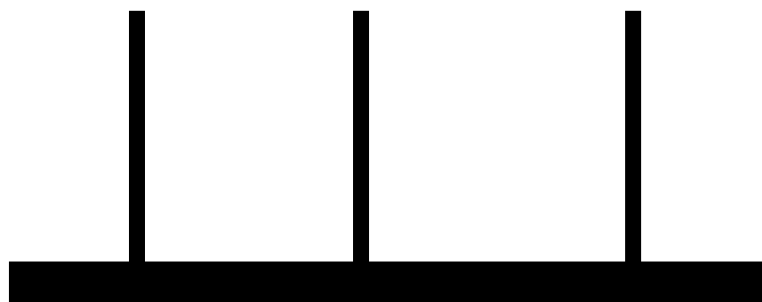
$$\mathbf{3} \otimes \mathbf{3} = \mathbf{6} \oplus \bar{\mathbf{3}}$$



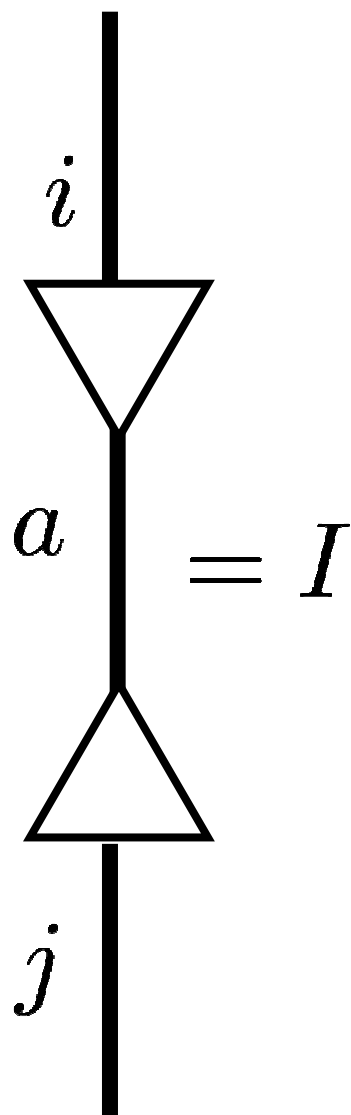
Case study 6. Graphical rep.



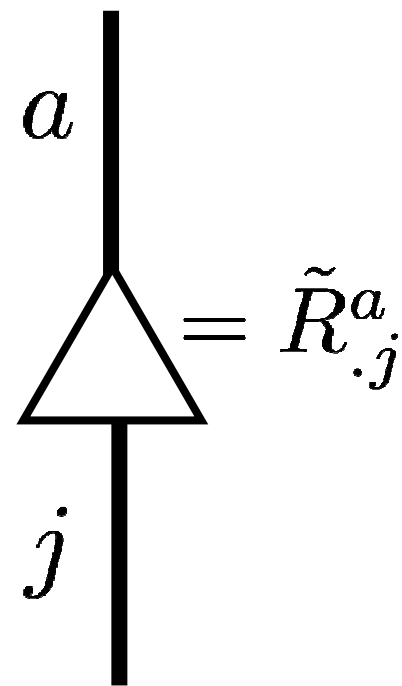
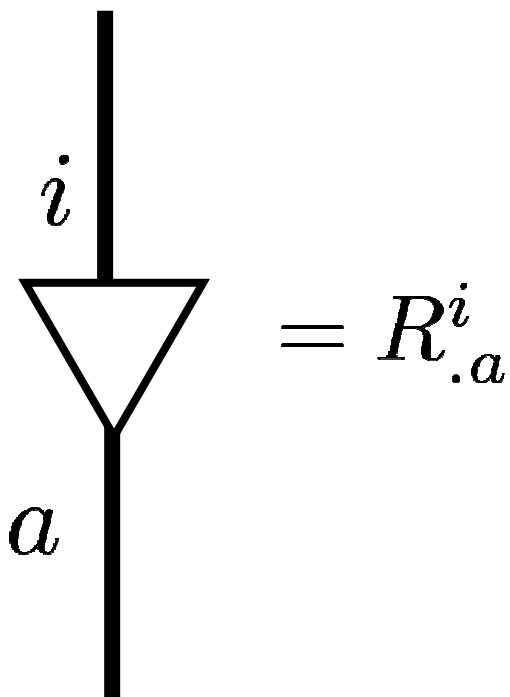
$$= \frac{1}{\sqrt{3!}} \epsilon_{ijk}$$



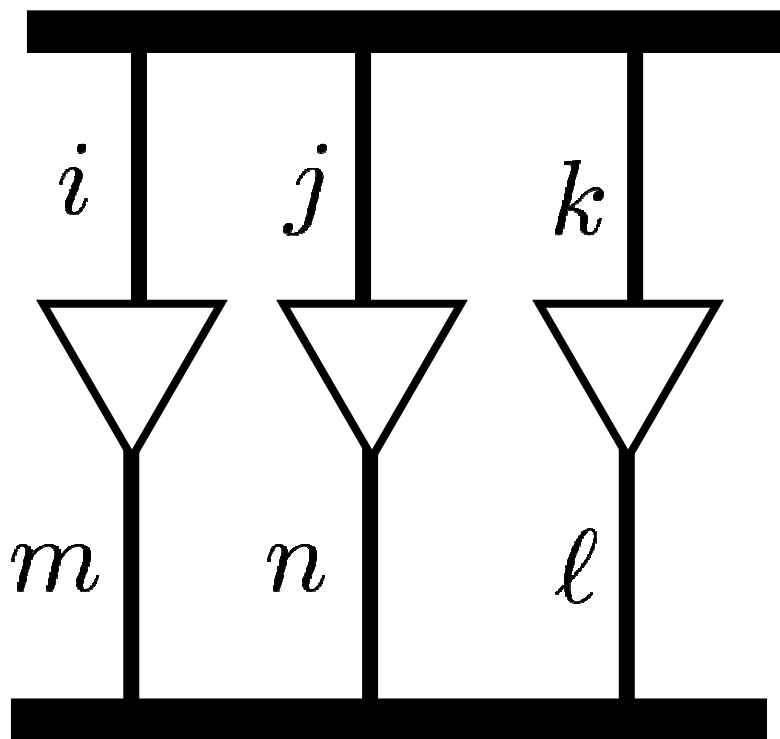
$$= \frac{1}{\sqrt{3!}} \epsilon^{ijk}$$



$$R\tilde{R} = 1, \quad R^i_{\cdot a} \tilde{R}^a_{\cdot j} = \delta^i_j$$

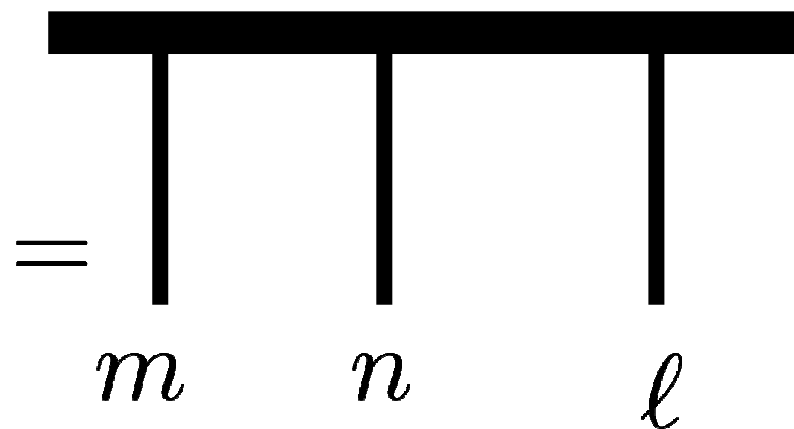
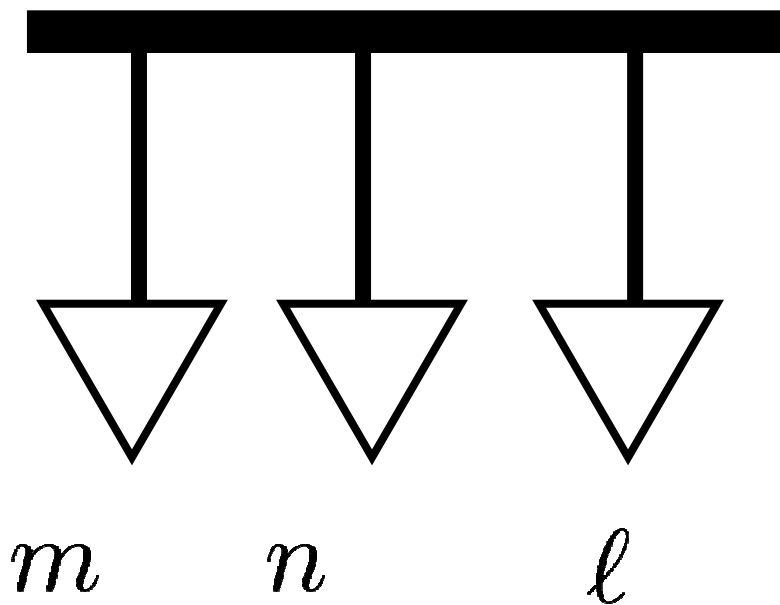


$$\det R = \frac{1}{3!} \epsilon_{ijk} \epsilon^{mnl} R^i_{.m} R^j_{.n} R^k_{.\ell} = 1$$

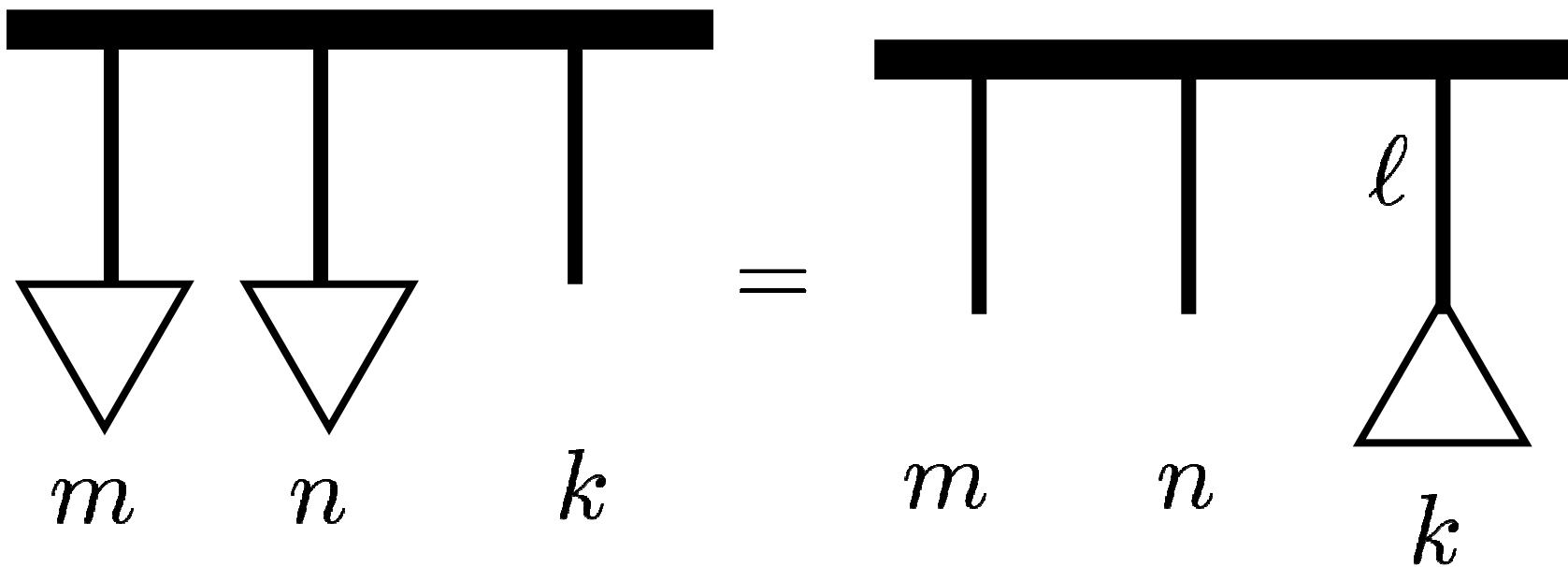


$$= 1$$

$$\epsilon_{ijk} R^i_{.m} R^j_{.n} R^k_{.\ell} = \epsilon_{mnl}$$



$$\epsilon_{ijk} R^i_{.m} R^j_{.n} = \epsilon_{mnl} R^{\cdot \ell}_k = \epsilon_{mnl} \tilde{R}^\ell_{.k}$$



Ref: Landau & Lifshitz, V2

Now consider four-tensors. The space components ($i, k, = 1, 2, 3$) of the antisymmetric tensor A^{ik} form a three-dimensional antisymmetric tensor with respect to purely spatial transformations; according to our statement its components can be expressed in terms of the components of a three-dimensional axial vector. With respect to these same transformations the components A^{01}, A^{02}, A^{03} form a three-dimensional polar vector. Thus the components of an antisymmetric four-tensor can be written as a matrix:

$$(A^{ik}) = \begin{vmatrix} 0 & p_x & p_y & p_z \\ -p_x & 0 & -a_z & a_y \\ -p_y & a_z & 0 & -a_x \\ -p_z & -a_y & a_x & 0 \end{vmatrix}, \quad (6.10)$$

where with respect to spatial transformations, $\mathbf{p}$ and $\mathbf{a}$ are polar and axial vectors, respectively.